GEOMETRIC HARDY INEQUALITIES ON STARSHAPED SETS

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ABSTRACT. In this paper, we present the geometric Hardy inequalities on the starshaped sets in the Carnot groups. Also, we obtain the geometric Hardy inequalities on half-spaces for general vector fields.

1. INTRODUCTION

In 1998, Danielli and Garofalo [4] firstly introduced the concept of starshapedness on the Carnot groups (see also [5]). Their paper provides the geometrical properties of starshaped and convex sets. The convexity in the Heisenberg groups was studied by many authors such as Monti and Rickly [10] who proved the geodesic convexity, or by Danielli, Garofalo, and Nhieu [3] (see also [8]) who introduced the concept of horizontal convexity (*H*-convexity). Bardi and Dragoni [1], [2] generalised the concept of convexity to general vector fields and introduced the notion of \mathcal{X} -convexity which is a generalisation of *H*-convexity. This analysis allows introducing the distance to the boundary notation for starshaped sets, so by using the distance formula one can obtain geometric Hardy type inequalities.

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The main aim of this paper is to obtain the geometric Hardy inequalities on starshaped sets in the Carnot groups. Moreover, we present the geometric Hardy inequalities on the half-spaces for general vector fields.

We organise the paper in the following way:

- Sec. 1: We give a brief overview of the sub-Riemannian manifolds, Grushin plane, Carnot groups, Heisenberg groups, and Engel groups.
- Sec. 2: We obtain the geometric Hardy inequalities on the starshaped sets in the Carnot groups and provide some examples.
- Sec. 3: We obtain the geometric Hardy inequalities on the half-spaces for general vector fields and provide some examples.
- Sec. 4: We give the proofs of main results.

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1.1. Sub-Riemannian manifolds. Let M be a smooth manifold of dimension n with a family of vector fields $\{X_k\}_{k=1}^N$, $n \geq N$, defined on M satisfying the Hörmander rank condition. Then they induce a sub-Riemannian metric $\langle \cdot, \cdot \rangle_{\mathcal{H}}$ on the associated space $\mathcal{H}_x = \operatorname{span}(X_1(x), \ldots, X_N(x))$. The triple $(M, \mathcal{H}, \langle \cdot, \cdot \rangle_{\mathcal{H}})$ is a so-called sub-Riemannian manifold (with sub-Riemannian geometry). Note that, unlike for Carnot groups, in general, it is not possible to define dilations, translations, the homogeneous norm and the distance on sub-Riemannian manifolds.

Let us denote the operator of the sum of squares of vector fields by

$$\mathcal{L} := \sum_{k=1}^{N} X_k^2. \tag{1.1}$$

These operators have been studied by many authors, for instance, it is well-known since Hörmander's pioneering work [9] that if the commutators of the vector fields $\{X_k\}_{k=1}^N$ generate the Lie algebra, the operator \mathcal{L} is locally hypoelliptic. The *p*-version of the sum of squares of vector fields can be given by the formula

$$\mathcal{L}_p f := \nabla_X \cdot (|\nabla_X f|^{p-2} \nabla_X f), \qquad (1.2)$$

where

$$\nabla_X := (X_1, \ldots, X_N).$$

1.2. Grushin plane. One of the important examples of a sub-Riemannian manifold is the Grushin plane. The Grushin plane is the space \mathbb{R}^2 with vector fields

$$X_1 = \frac{\partial}{\partial x_1}$$
, and $X_2 = x_1 \frac{\partial}{\partial x_2}$,

for $x := (x_1, x_2) \in \mathbb{R}^2$.

1.3. Carnot groups. Let $\mathbb{G} = (\mathbb{R}^n, \circ, \delta_{\lambda})$ be a stratified Lie group (or a homogeneous Carnot group or just a Carnot group), with the dilation structure δ_{λ} and Jacobian generators X_1, \ldots, X_N , so that N is the dimension of the first stratum of \mathbb{G} . Let us denote by Q the homogeneous dimension of \mathbb{G} . We refer to the recent books [7] and [15] for extensive discussions of stratified Lie groups and their properties.

The sub-Laplacian on \mathbb{G} is given by

$$\mathcal{L} = \sum_{k=1}^{N} X_k^2. \tag{1.3}$$

We also recall that the standard Lebesgue measure dx on \mathbb{R}^n is the Haar measure for \mathbb{G} (see, e.g. [7, Proposition 1.6.6]). Each left invariant vector field X_k has an explicit form and satisfies the divergence theorem, see e.g. [7] for the derivation of exact formula: more precisely, we can express

$$X_{k} = \frac{\partial}{\partial x'_{k}} + \sum_{l=2}^{r} \sum_{m=1}^{N_{l}} a_{k,m}^{(l)}(x', ..., x^{(l-1)}) \frac{\partial}{\partial x_{m}^{(l)}},$$
(1.4)

with $x = (x', x^{(2)}, \dots, x^{(r)})$, where r is the step of \mathbb{G} and $x^{(l)} = (x_1^{(l)}, \dots, x_{N_l}^{(l)})$ are the variables in the l^{th} stratum, see also [7, Section 3.1.5] for a general presentation. The

horizontal divergence is defined by

$$\operatorname{div}_H f := \nabla_H \cdot f,$$

where

$$\nabla_H := (X_1, \ldots, X_N)$$

is the horizontal gradient. The p-sub-Laplacian has the form

$$\mathcal{L}_p f = \nabla_H \cdot (|\nabla_H f|^{p-2} \nabla_H f).$$
(1.5)

1.4. Heisenberg groups. Let \mathbb{H}_1 be the Heisenberg group, that is, the set \mathbb{R}^3 equipped with the group law

$$x \circ x' := (x_1 + x'_1, x_2 + x'_2, x_3 + x'_3 + 2(x'_1x_2 - x_1x'_2)),$$

where $x := (x_1, x_2, x_3) \in \mathbb{R}^3$, and $x^{-1} = -x$ is the inverse element of x with respect to the group law. The dilation operation on the Heisenberg group with respect to the group law has the form

$$\delta_{\lambda}(x) := (\lambda x_1, \lambda x_2, \lambda^2 x_3) \text{ for } \lambda > 0.$$

The Lie algebra $\mathfrak h$ of the left-invariant vector fields on the Heisenberg group $\mathbb H_1$ is spanned by

$$X_1 := \frac{\partial}{\partial x_1} + 2x_2 \frac{\partial}{\partial x_3},$$
$$X_2 := \frac{\partial}{\partial x_2} - 2x_1 \frac{\partial}{\partial x_3},$$

with their (non-zero) commutator

$$[X_1, X_2] = -4\frac{\partial}{\partial x_3}.$$

The horizontal gradient on \mathbb{H}_1 is given by

$$\nabla_H := (X_1, X_2),$$

so the sub-Laplacian on \mathbb{H}_1 is given by

$$\mathcal{L} := X_1^2 + X_2^2.$$

The Heisenberg group is the most common example of a step 2 stratified group (Carnot group).

1.5. **Engel groups.** Let \mathbb{E} be the Engel group, that is, the set \mathbb{R}^4 equipped with the group law

$$x \circ x' := (x_1 + x'_1, x_2 + x'_2, x_3 + x'_3 + P_3, x_4 + x'_4 + P_4),$$

where

$$P_{3} = \frac{1}{2}(x_{1}x_{2}' - x_{2}x_{1}'),$$

$$P_{4} = \frac{1}{2}(x_{1}x_{3}' - x_{3}x_{1}') + \frac{1}{12}(x_{1}^{2}x_{2}' - x_{1}x_{1}'(x_{2} + x_{2}') + x_{2}x_{1}'^{2}).$$

Here $x := (x_1, x_2, x_3, x_4) \in \mathbb{R}^4$. The vector fields have the following form

$$X_{1} := \frac{\partial}{\partial x_{1}} - \frac{x_{2}}{2} \frac{\partial}{\partial x_{3}} - \left(\frac{x_{3}}{2} + \frac{x_{1}x_{2}}{12}\right) \frac{\partial}{\partial x_{4}},$$

$$X_{2} := \frac{\partial}{\partial x_{2}} + \frac{x_{1}}{2} \frac{\partial}{\partial x_{3}} + \frac{x_{1}^{2}}{12} \frac{\partial}{\partial x_{4}},$$

$$X_{3} := [X_{1}, X_{2}] = \frac{\partial}{\partial x_{3}} + \frac{x_{1}}{2} \frac{\partial}{\partial x_{4}},$$

$$X_{4} := [X_{1}, X_{3}] = \frac{\partial}{\partial x_{4}}.$$

The Engel group is a well-known example of a step 3 stratified group (Carnot group).

2. HARDY INEQUALITIES ON STARSHAPED SETS

In order to present the results on the starshaped domains, let us recall the definition of starshaped sets in a Carnot group $\mathbb{G} = (\mathbb{R}^n, \circ, \delta_t)$ and related arguments.

Definition 2.1 (Starshapedness [4]). Let $\Omega \subset \mathbb{G}$ be a C^1 domain containing the identity e. Then Ω is starshaped with respect to e if for every $x \in \partial \Omega$ one has

$$\langle Z(x), n(x) \rangle \ge 0, \tag{2.1}$$

where *n* is the Riemannian outer normal to $\partial \Omega$.

When the strict inequality holds, then Ω is said to be strictly starshaped with respect to e.

Here the vector fields Z are the infinitesimal generator of this group automorphism. This vector fields Z takes the form

$$Z = \sum_{i=1}^{N} x_i' \frac{\partial}{\partial x_i'} + 2 \sum_{l=1}^{N_2} x_{2,l} \frac{\partial}{\partial x_{2,l}} + \dots + r \sum_{l=1}^{N_r} x_{r,l} \frac{\partial}{\partial x_{r,l}}.$$
 (2.2)

Then for $x' \in \mathbb{R}^N$ and $x^{(i)} \in \mathbb{R}^{N_i}$ with $i = 2, \ldots, r$ we have

$$Z(x) = (x', 2x^{(2)}, \cdots, rx^{(r)}), \qquad (2.3)$$

and

$$\langle Z(x), n(x) \rangle = x'n' + 2x^{(2)}n^{(2)} + \ldots + rx^{(r)}n^{(r)}$$

= $x'_1n'_1 + \cdots + x'_Nn'_N + 2(x_{2,1}n_{2,1} + \cdots + x_{2,N_2}n_{2,N_2})$
+ $\cdots + r(x_{r,1}n_{r,1} + \cdots + x_{r,N_r}n_{r,N_r}),$

since $n(x) := (n', n^{(2)}, ..., n^{(r)})$ with $n' \in \mathbb{R}^N$ and $n^{(i)} \in \mathbb{R}^{N_i}, i = 2, ..., r$.

Based on the above arguments now we present the geometric Hardy inequalities on the starshaped sets for the sub-Laplacians. **Theorem 1.** Let Ω be a starshaped set on a Carnot group. Then for every $\gamma \in \mathbb{R}$ and p > 1 we have the following Hardy inequality

$$\int_{\Omega} |\nabla_H f(x)|^p dx \ge -(p-1)(|\gamma|^{\frac{p}{p-1}} + \gamma) \int_{\Omega} \frac{|\nabla_H \langle Z(x), n(x) \rangle|^p}{|\langle Z(x), n(x) \rangle|^p} |f(x)|^p dx \qquad (2.4)$$

$$+ \gamma \int_{\Omega} \frac{\mathcal{L}_p(\langle Z(x), n(x) \rangle)}{|\langle Z(x), n(x) \rangle|^{p-1}} |f(x)|^p dx,$$

for every function $f \in C_0^{\infty}(\Omega)$.

Corollary 2. Let \mathbb{H}^* be a starshaped set on the Heisenberg group \mathbb{H}_1 . Then for p > 1, we have the following Hardy inequality

$$\int_{\mathbb{H}^*} |\nabla_H f(x)|^p dx \ge \left(\frac{p-1}{p}\right)^p \int_{\mathbb{H}^*} \frac{|(n_1 + 4x_2n_3, n_2 - 4x_1n_3)|^p}{|x_1n_1 + x_2n_2 + 2x_3n_3|^p} |f(x)|^p dx, \quad (2.5)$$

for every function $f \in C_0^{\infty}(\mathbb{H}^*)$.

Remark 3. Note that in the case

$$\mathbb{H}^* := \{ \langle Z(x), n(x) \rangle > 0, \ \forall x \in \partial \mathbb{H}^*, \ Z(x) := (x_1, x_2, 2x_3) \} = \{ x \in \mathbb{H}_1 \cong \mathbb{R}^3 : x_3 > 0 \}$$

with n(x) := (0, 0, 1), and p = 2, we have the inequality

$$\int_{\mathbb{H}^*} |\nabla_H f(x)|^2 dx \ge \int_{\mathbb{H}^*} \frac{|x_1|^2 + |x_2|^2}{|x_3|^2} |f(x)|^2 dx.$$

Proof of Corollary 2. We begin the proof of Corollary 2 by a simple computation such as

$$\begin{aligned} \langle Z(x), n(x) \rangle &= x_1 n_1 + x_2 n_2 + 2x_3 n_3, \\ \nabla_H \langle Z(x), n(x) \rangle &= (n_1 + 4x_2 n_3, n_2 - 4x_1 n_3), \\ |\nabla_H \langle Z(x), n(x) \rangle|^p &= \left((n_1 + 4x_2 n_3)^2 + (n_2 - 4x_1 n_3)^2 \right)^{p/2}, \end{aligned}$$

and

$$\begin{aligned} \mathcal{L}_p \langle Z(x), n(x) \rangle = & \nabla_H \cdot (|\nabla_H \langle Z(x), n(x) \rangle|^{p-2} \nabla_H \langle Z(x), n(x) \rangle) \\ = & X_1 (|\nabla_H \langle Z(x), n(x) \rangle|^{p-2} (n_1 + 4x_2 n_3)) \\ & + & X_2 (|\nabla_H \langle Z(x), n(x) \rangle|^{p-2} (n_2 - 4x_1 n_3)) \\ = & - & 4(p-2) |\nabla_H \langle Z(x), n(x) \rangle|^{p-4} (n_1 + 4x_2 n_3) (n_2 - 4x_1 n_3) n_3 \\ & + & 4(p-2) |\nabla_H \langle Z(x), n(x) \rangle|^{p-4} (n_2 - 4x_1 n_3) (n_1 + 2x_4 n_3) n_3 \\ = & 0. \end{aligned}$$

Plugging the above expressions into inequality (2.4) and maximising with respect to γ , we arrive at inequality (2.5) which proves Corollary 2.

Corollary 4. Let \mathbb{E}^* be a starshaped set on the Engel group \mathbb{E} . Then for every function $f \in C_0^{\infty}(\mathbb{E}^*), \gamma \in \mathbb{R}$ and p = 2, we have

$$\int_{\mathbb{R}^*} |\nabla_H f(x)|^2 dx \ge - \left(|\gamma|^2 + \gamma\right) \int_{\mathbb{R}^*} \frac{|\nabla_H \langle Z(x), n(x) \rangle|^2}{\langle Z(x), n(x) \rangle^2} |f(x)|^2 dx \qquad (2.6)$$
$$+ \frac{\gamma}{2} \int_{\mathbb{R}^*} \frac{x_2 n_4}{\langle Z(x), n(x) \rangle} |f(x)|^2 dx.$$

Proof of Corollary 4. We begin the proof of Corollary 4 by a simple computation such as

$$\langle Z(x), n(x) \rangle = x_1 n_1 + x_2 n_2 + 2x_3 n_3 + 3x_4 n_4,$$

$$\nabla_H \langle Z(x), n(x) \rangle = \left(n_1 - x_2 n_3 - \frac{3x_3 n_4}{2} - \frac{x_1 x_2 n_4}{4}, n_2 + x_1 n_3 + \frac{x_1^2 n_4}{4} \right),$$

$$|\nabla_H \langle Z(x), n(x) \rangle|^2 = \left(n_1 - x_2 n_3 - \frac{3x_3 n_4}{2} - \frac{x_1 x_2 n_4}{4} \right)^2 + \left(n_2 + x_1 n_3 + \frac{x_1^2 n_4}{4} \right)^2,$$

and

$$\mathcal{L}\langle Z(x), n(x) \rangle = \nabla_H \cdot \nabla_H \langle Z(x), n(x) \rangle$$

= $X_1 \left(n_1 - x_2 n_3 - \frac{3x_3 n_4}{2} - \frac{x_1 x_2 n_4}{4} \right) + X_2 \left(n_2 + x_1 n_3 + \frac{x_1^2 n_4}{4} \right)$
= $\frac{x_2 n_4}{2}$.

Plugging the above expressions into inequality (2.4)

$$\int_{\mathbb{R}^*} |\nabla_H f(x)|^2 dx \ge - \left(|\gamma|^2 + \gamma\right) \int_{\mathbb{R}^*} \frac{|\nabla_H \langle Z(x), n(x) \rangle|^2}{\langle Z(x), n(x) \rangle^2} |f(x)|^2 dx + \frac{\gamma}{2} \int_{\mathbb{R}^*} \frac{x_2 n_4}{\langle Z(x), n(x) \rangle} |f(x)|^2 dx,$$

which proves Corollary 4.

3. Hardy inequalities on half-spaces for general vector fields

Let us define the half-space of a sub-Riemannian manifold by

$$\Omega^+ := \{ x \in \mathbb{R}^n : \langle x, n(x) \rangle > d \},\$$

where $n(x) \in \mathbb{R}^n$ is the Riemannian outer unit normal to $\partial \Omega^+$ and $d \in \mathbb{R}$. The Euclidean distance to the boundary $\partial \Omega^+$ is denoted by $dist(x, \partial \Omega^+)$ and defined by

$$dist(x,\partial\Omega^+) := \langle x, n(x) \rangle - d.$$

Then we have:

Theorem 5. Let M be a sub-Riemannian manifold, let $\Omega^+ \subset M$ be a half-space and let X_1, \ldots, X_N be the general vector fields. Then for every $\gamma \in \mathbb{R}$ and p > 1 we

have the following Hardy inequality

$$\int_{\Omega^{+}} |\nabla_{X} f(x)|^{p} dx \geq -(p-1)(|\gamma|^{\frac{p}{p-1}}+\gamma) \int_{\Omega^{+}} \frac{|\nabla_{X} dist(x,\partial\Omega^{+})|^{p}}{dist(x,\partial\Omega^{+})^{p}} |f(x)|^{p} dx \quad (3.1)$$
$$+ \gamma \int_{\Omega^{+}} \frac{\mathcal{L}_{p}(dist(x,\partial\Omega^{+}))}{dist(x,\partial\Omega^{+})^{p-1}} |f(x)|^{p} dx,$$

for every function $f \in C_0^{\infty}(\Omega^+)$.

Note that inequality (3.1) was obtained in the Carnot groups by the authors in [12], but here we extend it to general sub-Riemannian manifolds.

Let us give examples for the Heisenberg group (step 2), the Engel group (step 3), and the Grushin plane which does not have a group structure, but serves as an important example of the sub-Riemannian geometry.

Corollary 6. Let Ω^+ be a half-space in the Grushin plane G. Then for every function $f \in C_0^{\infty}(\Omega^+)$ and p > 1, we have the following Hardy inequality

$$\int_{\Omega^{+}} |\nabla_{X} f(x)|^{p} dx \geq -(p-1)(|\gamma|^{\frac{p}{p-1}}+\gamma) \int_{\Omega^{+}} \frac{(n_{1}^{2}+x_{1}^{2}n_{2}^{2})^{p/2}}{(x_{1}n_{1}+x_{2}n_{2}-d)^{p}} |f(x)|^{p} dx \quad (3.2)$$
$$+(p-2)\gamma \int_{\Omega^{+}} \frac{|\nabla_{X} dist(x,\partial\Omega^{+})|^{p-4}n_{1}n_{2}^{2}x_{1}}{(x_{1}n_{1}+x_{2}n_{2}-d)^{p-1}} |f(x)|^{p} dx.$$

If one of the cases n(x) = (1,0) or n(x) = (0,1) holds, then we have

$$\int_{\Omega^+} |\nabla_X f(x)|^p dx \ge \left(\frac{p-1}{p}\right)^p \int_{\Omega^+} \frac{(n_1^2 + x_1^2 n_2^2)^{p/2}}{(x_1 n_1 + x_2 n_2 - d)^p} |f(x)|^p dx,$$
(3.3)

where $dist(x, \partial \Omega^+) = \langle x, n(x) \rangle - d$ and $d \in \mathbb{R}$.

Remark 7. Note that, with ∇_X the Grushin gradient,

• If $\Omega^+ = \{x \in \mathbb{R}^2 : x_1 > d\}$ with n(x) = (1, 0), then we have

$$\int_{\Omega^+} |\nabla_X f(x)|^p dx \ge \left(\frac{p-1}{p}\right)^p \int_{\Omega^+} \frac{|f(x)|^p}{|x_1 - d|^p} dx$$

• If $\Omega^+ := \{x \in \mathbb{R}^2 : x_2 > d\}$ with n(x) = (0, 1), then we have

$$\int_{\Omega^+} |\nabla_X f(x)|^p dx \ge \left(\frac{p-1}{p}\right)^p \int_{\Omega^+} \frac{|x_1|^p}{|x_2 - d|^p} |f(x)|^p dx.$$

Proof of Corollary 6. We begin the proof of Corollary 6 by a simple computation such as

$$dist(x, \partial \Omega^{+}) = x_{1}n_{1} + x_{2}n_{2} - d,$$

$$\nabla_{X} dist(x, \partial \Omega^{+}) = (n_{1}, x_{1}n_{2}),$$

$$|\nabla_{X} dist(x, \partial \Omega^{+})|^{p} = (n_{1}^{2} + x_{1}^{2}n_{2}^{2})^{p/2},$$

and

$$\mathcal{L}_{p}dist(x,\partial\Omega^{+}) = \nabla_{X} \cdot (|\nabla_{X}dist(x,\partial\Omega^{+})|^{p-2} \nabla_{X}dist(x,\partial\Omega^{+}))$$

$$= \frac{\partial}{\partial x_{1}} ((n_{1}^{2} + x_{1}^{2}n_{2}^{2})^{\frac{p-2}{2}}n_{1}) + x_{1}\frac{\partial}{\partial x_{2}} ((n_{1}^{2} + x_{1}^{2}n_{2}^{2})^{\frac{p-2}{2}}x_{1}n_{2})$$

$$= (p-2)|\nabla_{X}dist(x,\partial\Omega^{+})|^{p-4}n_{1}n_{2}^{2}x_{1}.$$

Plugging the above expressions into inequality (3.1) we arrive at

$$\int_{\Omega^{+}} |\nabla_{X} f(x)|^{p} dx \ge -(p-1)(|\gamma|^{\frac{p}{p-1}}+\gamma) \int_{\Omega^{+}} \frac{(n_{1}^{2}+x_{1}^{2}n_{2}^{2})^{p/2}}{(x_{1}n_{1}+x_{2}n_{2}-d)^{p}} |f(x)|^{p} dx + (p-2)\gamma \int_{\Omega^{+}} \frac{|\nabla_{X} dist(x,\partial\Omega^{+})|^{p-4}n_{1}n_{2}^{2}x_{1}}{(x_{1}n_{1}+x_{2}n_{2}-d)^{p-1}} |f(x)|^{p} dx,$$

which proves inequality (3.2). If one of the cases n(x) = (1, 0) or n(x) = (0, 1) holds, then the last term of the above inequality vanishes, so that we get

$$\int_{\Omega^+} |\nabla_X f(x)|^p dx \ge -(p-1)(|\gamma|^{\frac{p}{p-1}} + \gamma) \int_{\Omega^+} \frac{(n_1^2 + x_1^2 n_2^2)^{p/2}}{(x_1 n_1 + x_2 n_2 - d)^p} |f(x)|^p dx.$$
(3.4)

Then, we maximise above inequality by differentiating with respect to γ , so that we have

$$\frac{p}{p-1}|\gamma|^{\frac{1}{p-1}} + 1 = 0,$$

which leads to

$$\gamma = -\left(\frac{p-1}{p}\right)^{p-1}.$$

By putting the value of γ into inequality (3.4), we obtain inequality (3.3).

Corollary 8. Let Ω^+ be a half-space on the Heisenberg group. Then for every function $f \in C_0^{\infty}(\Omega^+)$ and p > 1, we have

$$\int_{\Omega^+} |\nabla_H f(x)|^p dx \ge \left(\frac{p-1}{p}\right)^p \int_{\Omega^+} \frac{|(n_1 + 2x_2n_3, n_2 - 2x_1n_3)|^p}{dist(x, \partial\Omega^+)^p} |f(x)|^p dx, \quad (3.5)$$

where $dist(x, \partial\Omega^+) = \langle x, n(x) \rangle - d$ and $d \in \mathbb{R}$.

Remark 3.1. Note that if we choose n(x) = (0, 0, 1), p = 2 and d = 0 in inequality (3.5), then we get

$$\int_{\Omega^+} |\nabla_H f(x)|^2 dx \ge \int_{\Omega^+} \frac{|x_1|^2 + |x_2|^2}{|x_3|^2} |f(x)|^2 dx.$$
(3.6)

The Hardy inequality of the form (3.6) in the half-space on the Heisenberg group was shown by Luan and Young [16].

Proof of Corollary 8. We begin the proof of Corollary 8 by a simple computation such as

$$dist(x,\partial\Omega^{+}) = x_{1}n_{1} + x_{2}n_{2} + x_{3}n_{3} - d,$$

$$\nabla_{X}dist(x,\partial\Omega^{+}) = (n_{1} + 2x_{2}n_{3}, n_{2} - 2x_{1}n_{3}),$$

$$|\nabla_{X}dist(x,\partial\Omega^{+})|^{p} = ((n_{1} + 2x_{2}n_{3})^{2} + (n_{2} - 2x_{1}n_{3})^{2})^{p/2}.$$

Then we compute

$$\begin{aligned} \mathcal{L}_{p}dist(x,\partial\Omega^{+}) = &\nabla_{H} \cdot (|\nabla_{H}dist(x,\partial\Omega^{+})|^{p-2}\nabla_{H}dist(x,\partial\Omega^{+})) \\ = &X_{1}((n_{1}+2x_{2}n_{3})^{2}+(n_{2}-2x_{1}n_{3})^{2})^{\frac{p-2}{2}}(n_{1}+2x_{2}n_{3}) \\ &+X_{2}((n_{1}+2x_{2}n_{3})^{2}+(n_{2}-2x_{1}n_{3})^{2})^{\frac{p-2}{2}}(n_{2}-2x_{1}n_{3}) \\ = &-2(p-2)|\nabla_{H}dist(x,\partial\Omega^{+})|^{p-4}(n_{1}+2x_{2}n_{3})(n_{2}-2x_{1}n_{3})n_{3} \\ &+2(p-2)|\nabla_{H}dist(x,\partial\Omega^{+})|^{p-4}(n_{2}-2x_{1}n_{3})(n_{1}+2x_{2}n_{3})n_{3} \\ =&0. \end{aligned}$$

Plugging the above expressions into inequality (3.1), we arrive at

$$\int_{\Omega^{+}} |\nabla_{H}f(x)|^{p} dx \ge -(p-1)(|\gamma|^{\frac{p}{p-1}} + \gamma) \int_{\Omega^{+}} \frac{|(n_{1} + 2x_{2}n_{3}, n_{2} - 2x_{1}n_{3})|^{p}}{(x_{1}n_{1} + x_{2}n_{2} + x_{3}n_{3} - d)^{p}} |f(x)|^{p} dx,$$
(3.7)

which can be maximised by differentiating with respect to γ , then we have

$$\frac{p}{p-1}|\gamma|^{\frac{1}{p-1}} + 1 = 0,$$

that leads to

$$\gamma = -\left(\frac{p-1}{p}\right)^{p-1}.$$

By putting the value of γ into inequality (3.7), we obtain inequality

$$\int_{\Omega^+} |\nabla_H f(x)|^p dx \ge \left(\frac{p-1}{p}\right)^p \int_{\Omega^+} \frac{|(n_1 + 2x_2n_3, n_2 - 2x_1n_3)|^p}{dist(x, \partial\Omega^+)^p} |f(x)|^p dx,$$
proves Corollary 8.

which proves Corollary 8.

Corollary 9. Let Ω^+ be a half-space on the Engel group \mathbb{E} . Then for every function $f \in C_0^{\infty}(\Omega^+)$, $\gamma \in \mathbb{R}$ and p = 2, we have

$$\int_{\Omega^{+}} |\nabla_{H}f(x)|^{2} dx \geq -\left(|\gamma|^{2}+\gamma\right) \int_{\Omega^{+}} \frac{|\nabla_{H}dist(x,\partial\Omega^{+})|^{2}}{dist(x,\partial\Omega^{+})^{2}} |f(x)|^{2} dx \qquad (3.8)$$
$$+ \frac{\gamma}{6} \int_{\Omega^{+}} \frac{x_{2}n_{4}}{dist(x,\partial\Omega^{+})} |f(x)|^{2} dx,$$

where $dist(x, \partial \Omega^+) = \langle x, n(x) \rangle - d$ and $d \in \mathbb{R}$.

Proof of Corollary 9. We begin the proof of Corollary 9 by a simple computation such as

$$dist(x,\partial\Omega^{+}) = x_{1}n_{1} + x_{2}n_{2} + x_{3}n_{3} + x_{4}n_{4} - d,$$

$$\nabla_{H}dist(x,\partial\Omega^{+}) = \left(n_{1} - \frac{x_{2}n_{3}}{2} - \frac{x_{3}n_{4}}{2} - \frac{x_{1}x_{2}n_{4}}{12}, n_{2} + \frac{x_{1}n_{3}}{2} + \frac{x_{1}^{2}n_{4}}{12}\right),$$

$$\nabla_{H}dist(x,\partial\Omega^{+})|^{2} = \left(n_{1} - \frac{x_{2}n_{3}}{2} - \frac{x_{3}n_{4}}{2} - \frac{x_{1}x_{2}n_{4}}{12}\right)^{2} + \left(n_{2} + \frac{x_{1}n_{3}}{2} + \frac{x_{1}^{2}n_{4}}{12}\right)^{2},$$

and

$$\begin{aligned} \mathcal{L}(dist(x,\partial\Omega^{+})) = &\nabla_{H} \cdot \nabla_{H} dist(x,\partial\Omega^{+}) \\ = &\nabla_{H} \cdot \left(n_{1} - \frac{x_{2}n_{3}}{2} - \frac{x_{3}n_{4}}{2} - \frac{x_{1}x_{2}n_{4}}{12}, n_{2} + \frac{x_{1}n_{3}}{2} + \frac{x_{1}^{2}n_{4}}{12} \right) \\ = &X_{1} \left(n_{1} - \frac{x_{2}n_{3}}{2} - \frac{x_{3}n_{4}}{2} - \frac{x_{1}x_{2}n_{4}}{12} \right) + X_{2} \left(n_{2} + \frac{x_{1}n_{3}}{2} + \frac{x_{1}^{2}n_{4}}{12} \right) \\ = &\frac{x_{2}n_{4}}{6}. \end{aligned}$$

Plugging the above expressions into inequality (3.1), we get

$$\int_{\Omega^{+}} |\nabla_{H}f(x)|^{2} dx \geq -(|\gamma|^{2}+\gamma) \int_{\Omega^{+}} \frac{|\nabla_{H}dist(x,\partial\Omega^{+})|^{2}}{dist(x,\partial\Omega^{+})^{2}} |f(x)|^{2} dx$$
$$+ \frac{\gamma}{6} \int_{\Omega^{+}} \frac{x_{2}n_{4}}{dist(x,\partial\Omega^{+})} |f(x)|^{2} dx,$$
 by es Corollary 9.

which proves Corollary 9.

4. PROOF OF MAIN RESULTS

The approach to prove the main results is based on the works [11] and [12] (see, also [13]-[14]). For a vector field $g \in C^{\infty}(\Omega)$ we compute

$$\begin{split} \int_{\Omega} \operatorname{div}_X g |f(x)|^p dx &= -p \int_{\Omega} |f(x)|^{p-1} \langle g, \nabla_X f(x) \rangle dx \\ &\leq p \left(\int_{\Omega} |\nabla_X f(x)|^p dx \right)^{\frac{1}{p}} \left(\int_{\Omega} |g|^{\frac{p}{p-1}} |f(x)|^p dx \right)^{\frac{p-1}{p}} \\ &\leq \int_{\Omega} |\nabla_H f(x)|^p dx + (p-1) \int_{\Omega} |g|^{\frac{p}{p-1}} |f(x)|^p dx. \end{split}$$

Here we have first used the divergence theorem, then we applied the Hölder inequality and the Young inequality. By rearranging the above expression, we arrive at

$$\int_{\Omega} |\nabla_X f(x)|^p dx \ge \int_{\Omega} (\operatorname{div}_X g - (p-1)|g|^{\frac{p}{p-1}})|f(x)|^p dx.$$
(4.1)

A suitable choice of the vector field g in each special case is a key argument of our proofs.

Proof of Theorem 1. Let us set

$$g = \gamma \frac{|\nabla_H \langle Z(x), n(x) \rangle|^{p-2}}{|\langle Z(x), n(x) \rangle|^{p-1}} \nabla_H \langle Z(x), n(x) \rangle,$$

so that we have

$$|g|^{\frac{p}{p-1}} = |\gamma|^{\frac{p}{p-1}} \frac{|\nabla_H \langle Z(x), n(x) \rangle|^p}{|\langle Z(x), n(x) \rangle|^p},$$
(4.2)

and

$$\operatorname{div}_{H}g = \gamma \frac{\mathcal{L}_{p}(\langle Z(x), n(x) \rangle)}{|\langle Z(x), n(x) \rangle|^{p-1}} - \gamma(p-1) \frac{|\nabla_{H} \langle Z(x), n(x) \rangle|^{p}}{|\langle Z(x), n(x) \rangle|^{p}}.$$
(4.3)

Plugging the above expressions (4.2) and (4.3) into inequality (4.1), we get

$$\int_{\Omega} |\nabla_H f(x)|^p dx \ge -(p-1)(|\gamma|^{\frac{p}{p-1}} + \gamma) \int_{\Omega} \frac{|\nabla_H \langle Z(x), n(x) \rangle|^p}{|\langle Z(x), n(x) \rangle|^p} |f(x)|^p dx + \gamma \int_{\Omega} \frac{\mathcal{L}_p(\langle Z(x), n(x) \rangle)}{|\langle Z(x), n(x) \rangle|^{p-1}} |f(x)|^p dx,$$

which proves inequality (2.4).

Proof of Theorem 5. Let us take

$$g = \gamma \frac{|\nabla_X dist(x, \partial \Omega^+)|^{p-2}}{dist(x, \partial \Omega^+)^{p-1}} \nabla_X dist(x, \partial \Omega^+), \qquad (4.4)$$

so that we have

$$|g|^{\frac{p}{p-1}} = |\gamma|^{\frac{p}{p-1}} \frac{|\nabla_X dist(x, \partial\Omega^+)|^p}{dist(x, \partial\Omega^+)^p},\tag{4.5}$$

and

$$\operatorname{div}_X g = \gamma \frac{\mathcal{L}_p \operatorname{dist}(x, \partial \Omega^+)}{\operatorname{dist}(x, \partial \Omega^+)^{p-1}} - \gamma (p-1) \frac{|\nabla_X \operatorname{dist}(x, \partial \Omega^+)|^p}{\operatorname{dist}(x, \partial \Omega^+)^p}.$$
(4.6)

Combining expressions (4.5) and (4.6) with inequality (4.1), we obtain

$$\int_{\Omega} |\nabla_X f(x)|^p dx \ge -(p-1)(|\gamma|^{\frac{p}{p-1}} + \gamma) \int_{\Omega} \frac{|\nabla_X dist(x, \partial\Omega^+)|^p}{dist(x, \partial\Omega^+)^p} |f(x)|^p dx + \gamma \int_{\Omega} \frac{\mathcal{L}_p dist(x, \partial\Omega^+)}{dist(x, \partial\Omega^+)^{p-1}} |f(x)|^p dx,$$

which proves inequality (3.1).

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