

Bright and efficient sensitized near-infrared photoluminescence from an organic neodymium-containing composite material system.

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Supporting Information Placeholder

ABSTRACT: Intense organic neodymium (Nd^{3+}) emission is obtained with the near-infrared (NIR) emission equivalent in intensity to an organic semiconductor emitting material. The advantage of the Nd^{3+} emission is its narrow linewidth and NIR emission, which is enhanced by ~ 3000 times at low excitation power through an efficient sensitization effect from a composite organic sensitizer. This performance is optimized at high concentrations of Nd^{3+} ions and the organic perfluorinated system provides the ion excitations with a quantum efficiency of $\sim 40\%$. The material system is applicable to thin films that are compatible with integrated optics applications.

Near-infrared (NIR) emitting material systems are advantageous to fabricate light sources for bioimaging devices, optical sensors and optical communications¹⁻⁶, as NIR wavelengths are less attenuated than visible light in these applications, fitting important optical wavelengths including the second bioimaging window and the o-band optical communications window. Given a suitable material candidate for these applications, high luminescence intensity is of vital importance and this performance will rely on the absorption/emission cross-sections ($\sigma_{\text{abs}}/\sigma_{\text{em}}$), radiative rates (R_{rad}), excited-state densities (N_{ex}), and quantum efficiencies (η). For example, semiconductor-like nanomaterials such as organic fluorescent compounds, quantum dots, perovskite, etc⁷⁻¹² can give intense visible emission (usually $\ll 900$ nm) with high quantum efficiencies due to their large radiative rates ($R_{\text{rad}} = \sim 10^8$ to 10^9 s⁻¹) and large cross-sections ($\sigma_{\text{abs}}/\sigma_{\text{em}} = \sim 10^{-15}$ to 10^{-17} cm²). Nevertheless, tuning the emission wavelength deeper into the NIR region usually fundamentally changes their chemistry,^{13,14} which detrimentally affects performance. NIR lanthanide materials such as neodymium-, ytterbium- and erbium-doped materials utilize the advantages of their intrinsic NIR emission at > 1 μm , therefore lanthanide doped nanoparticles or organic complexes have long been under development as alternatives.^{15,16} However, the drawback is that the partially forbidden 4f-4f transitions of lanthanide ions result in weak emission/absorption cross-section ($\sigma_{\text{em}}/\sigma_{\text{abs}} = \sim 10^{-21}$ to 10^{-20} cm²) and a low radiative rate ($R_{\text{rad}} = \sim 10^2$ to 10^3 s⁻¹). These parameters are orders of magnitude smaller than in semiconductor materials and suggest very high excitation powers are inevitably required to produce strong luminescence. Also, high lanthanide ion concentrations induce a significant self-quenching effect^{17,18} that tends to limit the intensity that can be achieved. Moreover, NIR lanthanide containing materials are easily quenched by hydrogen containing organic moieties such as C-H, O-H, N-H, etc.^{19,20} This quenching effect is attributed to vibrational modes of those chemical bonds and can

greatly reduce the quantum efficiency of the NIR emission down to $< 0.1\%$ ^{20,21}.

In this work, we report an exceptionally intense organic-semiconductor/neodymium (Nd^{3+}) material system emitting at 1.06 μm and 1.3 μm utilising a molecular composite of an organic neodymium (Nd^{3+}) molecule and one organic chromophore semiconductor molecule. The Nd^{3+} ion is coordinated within a cage of perfluorinated ligands to eliminate vibrational quenching so that its internal quantum efficiency (IQY) can reach $\sim 40\%$, and the composite organic semiconductor chromophore gives an enormous sensitization effect to enhance the Nd^{3+} emission. The intense 1.06 μm emission fits the second bioimaging window; the 1.3 μm emission matches the optic-communication O-band, and the material system is easy to fabricate as thin films. These properties suggest a potential candidate of integrated light sources for those applications. The perfluorinated Nd^{3+} complex is $\text{Nd}(\text{F-TPIP})_3$ where the Nd^{3+} ion is incorporated with three tetrakis(pentafluorophenyl)imidodiphosphate (F-TPIP^-) ligands. $\text{Zn}(\text{F-BTZ})_2$ ^{22,23}, a perfluorinated analog of $\text{Zn}(\text{BTZ})_2$, which is an organic semiconductor material used for fabricating OLEDs, is the chromophore sensitizer component, where F-BTZ^- refers to the conjugate base of 2-(tetrafluoro-2-hydroxyphenyl)tetrafluorobenzothiazole. The enormous enhancement of Nd^{3+} emission by a factor of ~ 3000 is attributed to the energy transfer (ET) from the $\text{Zn}(\text{F-BTZ})_2$ excitons to the Nd^{3+} ions. It is noteworthy that the luminescence intensity of Nd^{3+} at the NIR range is almost as high as the emission intensity of $\text{Zn}(\text{F-BTZ})_2$ at the visible. That gives good evidence of intense NIR luminescence from Nd^{3+} doped material system with a low pump power and reasonably high Nd^{3+} concentration.

Result

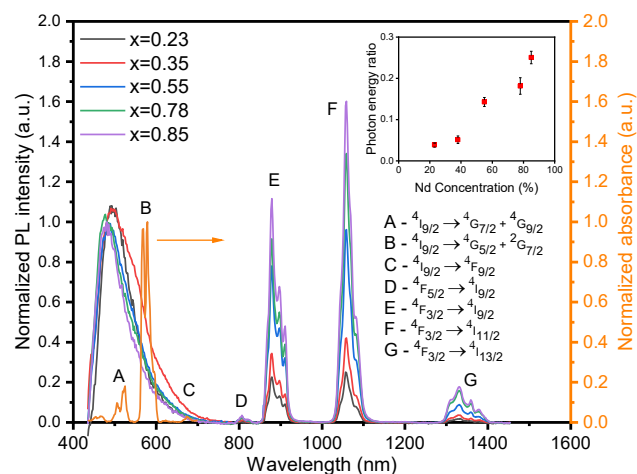


Figure 1. Absorption spectrum of Nd(F-TPIP)₃ and PL spectrum of [Nd(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} ($x = 0.23, 0.38, 0.55, 0.78$ and 0.85) thin films. The thickness of Zn(F-BTZ)₂ in these thin films is fixed to 50 nm. The excitation wavelength for PL spectra is 405 nm. All of the PL spectra are corrected for the responsivity of the utilized photon detectors and spectrometer gratings. The insert indicates the emission photon energy ratio between Nd³⁺ and Zn(F-BTZ)₂.

The Nd³⁺ absorption spectrum is obtained on an Nd(F-TPIP)₃ crystal that was grown by slow evaporation, shown as the orange curve in Figure 1 (to the left axis). This absorption spectrum demonstrates the characteristic Nd³⁺ transitions in the visible range and shows good overlap with the Zn(F-BTZ)₂ emission band. It is noteworthy that the absorption peak centred at 570 nm is a hypersensitive transition (⁴I_{9/2} → ⁴G_{5/2} + ⁴G_{7/2}) and its transition strength is 5 to 10 times stronger than the others.

A series of composite thin films, [Nd(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} ($x = 0.23, 0.38, 0.55, 0.78, 0.85$) in which Zn(F-BTZ)₂ is employed at varied concentrations, were fabricated to characterise the sensitized

Nd³⁺ photoluminescence (PL). A 405 nm diode laser photoexcites these films to give the sensitized PL spectra in Figure 1 (to the left axis), corresponding to the intrinsic Nd³⁺ transitions as well as a broad emission band from 400 nm to 750 nm due to the Zn(F-BTZ)₂ emission. The Nd³⁺ PL intensity increases at higher Nd(F-TPIP)₃ concentrations and the Nd³⁺ 1060 nm peak emission becomes higher than that of the Zn(F-BTZ)₂ emission. Integrating all the Nd³⁺ emission bands and the Zn(F-BTZ)₂ band on an energy scale results in the ratios of the emitted energy of the Nd³⁺ compared to that from Zn(F-BTZ)₂, shown in the inset in Figure 1. As the Nd(F-TPIP)₃ concentration increases, the ratio rises from ~0.04 to ~0.27 and demonstrates that up to 27% of the total energy emitted is in the NIR region due to the Nd³⁺. It is noteworthy that the Zn(F-BTZ)₂ emission cross-section is ~10⁻¹⁶ cm² with a nearly unity quantum efficiency for the fluorescence at room temperature (RT),²² which is 3 orders of magnitude larger than the Nd³⁺ intrinsic absorption/emission cross-section of ~10⁻¹⁹ cm². Hence, the similar intensity implies the chromophore has generated a significant population of Nd³⁺ excited states.

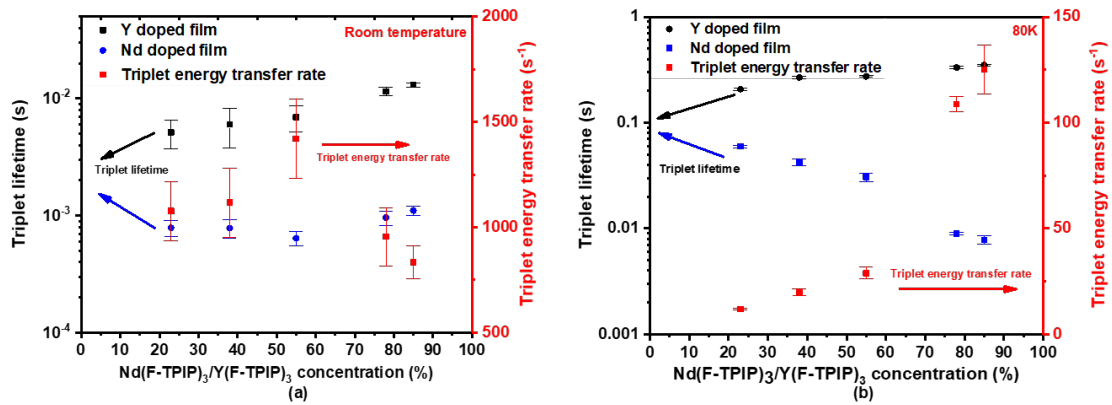


Figure 2. 550nm lifetime of [Nd(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} and [Y(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} thin films ($x = 0.23, 0.38, 0.55, 0.78$ and 0.85) at room temperature and 80K. Black dots indicate triplet lifetime in Y³⁺ doped film, blue dots represent triplet lifetime in Nd³⁺ doped films. The triplet energy transfer rate is given by red squares.

Time-resolved photoluminescence (TRPL) spectra of the composite films were recorded at 550 nm where the Zn(F-BTZ)₂ phosphorescence predominates due to the triplet emission²⁴. A series of composite films of [Y(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} ($x = 0.23, 0.38, 0.55, 0.78, 0.85$) were used to compare the triplet emission properties by excluding the ET as Y³⁺ ions are optically inert. These decay traces are fitted with a three exponential decay function to give the lifetimes of triplet states, shown in the Table S1 to S4. In Figure 2 (a), the triplet lifetimes clearly become shorter in the presence of Nd(F-TPIP)₃ compared to the triplet lifetimes in the Y³⁺ doped films at room temperature. This indicates Nd³⁺ ions significantly quench triplet excitons due to the sensitization. The difference in the triplet lifetimes can be used to calculate the triplet ET rates (~800 to ~1300 s⁻¹) to Nd³⁺ excitations in Table S1 and S2. Also,

we calculate the Förster energy transfer rate (FRET) between Zn(F-BTZ)₂ and Nd³⁺ to give an estimated triplet ET rate varied from 728 ± 213 s⁻¹ to 6186 ± 1814 s⁻¹, which are of the same order as the measured triplet ET rates, detailed in the SI. With the samples cooled to 80K, the thermally induced nonradiative quenching is reduced so that the triplet ET process to Nd³⁺ ions becomes more obvious where the triplet lifetimes can be quenched to ~7 ms from those that can be as long as ~330 ms in the Y(F-TPIP)₃ present samples. More evidently, a higher Nd(F-TPIP)₃ concentration makes the triplet lifetime much shorter and decreases from 40 ms to 7 ms. This suggests the triplet ET process is improved at a higher Nd³⁺ concentration, which might be expected as there are more Nd³⁺ ions to interact with each triplet.

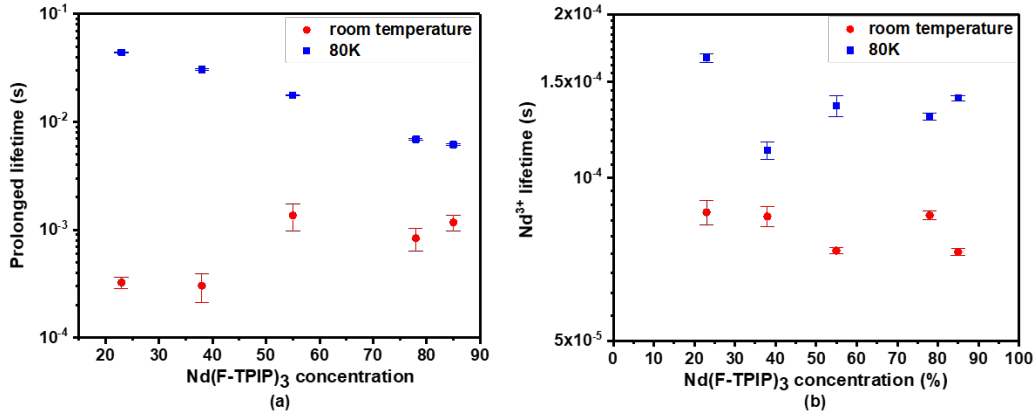


Figure 3. 1060 nm lifetime of [Nd(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} thin films ($x = 0.23, 0.38, 0.55, 0.78$ and 0.85) at room temperature and 80K. (a) Prolonged lifetime. (b) Nd³⁺ lifetime.

In Figure 3, the lifetime of Nd³⁺ 1060 nm is obtained to study of the influence on the Nd³⁺ depopulation process via the sensitization by Zn(F-BTZ)₂. At RT, they are around 80 μs, which correspond to a Nd³⁺ intrinsic decay process with an internal quantum efficiency (IQE) of $\sim 24.7 \pm 2.1 \%$, if we use an average literature reported value of 325 μs as the Nd³⁺ radiative lifetime^{25,26}. There is another decay process with weak emission intensity but longer lifetime (blue dots), from 300 μs to 1.3 ms at different concentration films. Although the fitting of these long lifetimes has relatively large error bars, these prolonged lifetimes are probably longer than the radiative lifetime. However, these prolonged lifetimes cannot be seen in the TRPL spectrum of neat Nd(F-TPIP)₃ powder (Figure S2). These prolonged lifetimes have as comparable values as the triplet lifetimes in those Nd(F-TPIP)₃ present films according to Figure 2. Thus, we attribute them to the ET process from triplet as the long-

lived triplet states, longer than the Nd³⁺ excited states, likely for them to re-excite some de-activated Nd³⁺ ions repeatedly. This process enables the replenishment of the excited state of Nd³⁺ ions until all the triplet states decay to the ground state. We believe this mechanism is consistent with those what have been presented and studied in a previous study of sensitized Yb³⁺ PL by Zn(F-BTZ)₂²⁷. At 80K, the prolonged lifetimes are significantly extended as shown in Figure 3 (a) and (b). These prolonged lifetimes and their concentration dependence are again consistent with the measured triplet lifetimes in the Nd(F-TPIP)₃ composited films at 80K. Furthermore, the intrinsic lifetimes of Nd³⁺ 1060 nm PL are extended to $\sim 125 \mu\text{s}$, which improves the IQE of Nd³⁺ ions to $38.7 \pm 5.3 \%$. That implies that low temperature restricts some thermally related non-radiative quenching routes to Nd³⁺ excited states.

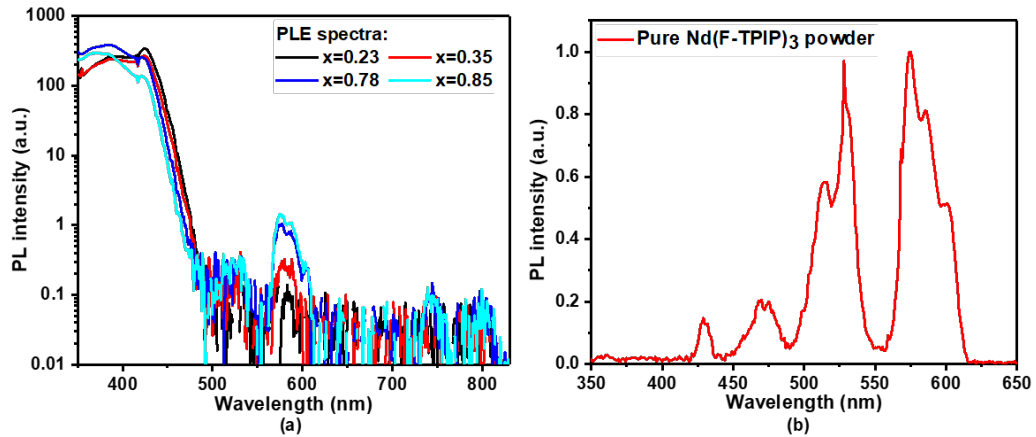


Figure 4. (a) PLE spectrum of [Nd(F-TPIP)₃]_x[Zn(F-BTZ)₂]_{1-x} thin films ($x = 0.23, 0.38, 0.78$ and 0.85). (b) PLE spectrum of pure Nd(F-TPIP)₃ powder.

Figure 4a shows the photoluminescence excitation (PLE) spectra of different Nd(F-TPIP)₃ concentration present films at a wavelength of 1060 nm to visualize the sensitization from Zn(F-BTZ)₂. Those weak PLE peaks (at 520 nm, 570nm, 780 nm and 808 nm) are due to the intrinsic Nd³⁺ absorption. The PLE band from 350 nm to 450 nm corresponds to the absorption range of Zn(F-BTZ)₂ which is shown in Figure S1. The PLE spectrum (Figure 4b) of a pure Nd(F-TPIP)₃ powder shows only the direct excitation into the intrinsic Nd³⁺ transitions without a F-TPIP⁻ contribution. Therefore, the ratio between the PL intensity generated by the absorption wavelength of Zn(F-BTZ)₂ and Nd³⁺ is a reliable measurement of the PL intensity enhancement via sensitization.

Figure S11 gives the PLE intensity enhancement with the Nd(F-TPIP)₃ concentration dependence. The enhancement factor is

obtained from the PLE peak intensity ratio around 405 nm and 570 nm, which means excitation indirectly into Zn(F-BTZ)₂ and directly Nd³⁺ ions. The enhancement factor increases the sensitization from 400 for the 80% Nd(F-TPIP)₃ to >3000 for the 20% Nd(F-TPIP)₃. Rate equations are used to simulate the dynamics of the excited population of Nd³⁺ ions, elaborated in the SI. Equation S11 results in the simulated enhancement factor that is equivalent to the ratio of the sensitized Nd³⁺ population ($N_{Nd1_sensitize}/N_{Nd1_direct}$) and shows the simulated curve is consistent with our experimentally measured data with an estimated singlet ET rate of $8 \times 10^8 \text{ s}^{-1}$. In the modelling, the photon-emission rate of the excited Nd³⁺ ions ($R_{Nd} \times 24.7\% \text{ IQE} = \sim 2.8 \times 10^3 \text{ Hz}$) is about 5 orders of magnitude lower than that of the singlet emission ($R_s = \sim 5 \times 10^8 \text{ Hz}$). However, the sensitization effect increases the Nd³⁺ excitation densities ($N_{Nd1_sensitize}$) by orders of magnitude to give a photon flux (\sim

$N_{Nd1_sensitize} \times R_{Nd}$) comparable to that of the singlet emission ($\sim N_s \times R_s$), which results in the comparable emission intensities of the Zn(F-BTZ)₃ fluorescence and the Nd(F-TPIP)₃ emission displayed in Figure 1.

In conclusion, efficient sensitization from an organic sensitizer Zn(F-BTZ)₂ to Nd³⁺ ions enhance the excitation densities of Nd³⁺ ions by ~ 3000 times maximally at low excitation power. The enhancement results in comparable Nd³⁺ PL intensities to the fluorescence of the chromophore. The excited state dynamics are investigated by rate equation simulation to explain the concentration dependence of the phenomenon. Our results suggest a potential NIR emitting organic material system with high power efficiencies for luminescence at 1060nm and 1330 nm.

ASSOCIATED CONTENT

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NOTES

The authors declare no competing financial interests.

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SUPPORTING INFORMATION

Experimental details and supporting data, including Figure S1- S13 and Table S1 – S12.

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TOC Graphic

