Room-temperature multiferroic behaviour in layer structured Aurivillius

phase ceramics

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ABSTRACT

Multiferroics that simultaneously exhibit ferroelectricity (FE) and ferromagnetism (FM) have recently attracted great attention due to their potential application in next generation electronic devices. However, only a few single-phase multiferroic materials exhibit ferroelectric and ferromagnetic order at room temperature. Recently, some bismuth layer-structured Aurivillius compounds were reported as multiferroics at room temperature, but the origin of their magnetic property is still under debate because the net magnetization may originate from the presence of secondary phases which are not easily detected by laboratory XRD diffractometers. Here, textured Aurivillius phase $Bi_{5.25}La_{0.75}FeCoTi_3O_{18}$ ceramics were prepared by Spark Plasma Sintering. The ferromagnetic character of the ceramics was indicated by the magnetic field-induced reversible intensity changes of a certain set of crystalline planes belonging to the Aurivillius phase, as measured by in-situ neutron diffraction under applied magnetic field. The first principle calculations indicate that the ferromagnetism originates from double exchange interactions Fe^{3+}-O-Fe^{3+}, Co^{3+}-O-Co^{3+} and Fe^{3+}-O-Co^{3+} in the ferro-toroidal main phase. Magnetic-controlled ferroelectric domain switching was observed

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by piezoelectric force microscopy at room temperature. The prepared Aurivillius phase ceramics, with Co/Fe contributing to magnetization and polarization at the same time, can be considered an intrinsic room-temperature multiferroic.

Multiferroics that simultaneously exhibit ferroelectricity (FE) and ferromagnetism (FM) have recently attracted great attention due to their potential application in new types of sensors, actuators, memory storage devices, etc.[1-4] However, only a few single-phase multiferroic materials exhibit ferroelectric and ferromagnetic order at room temperature.[5] Recently, bismuth layer-structured Aurivillius compounds of a general formula $(Bi_2O_2)^{2+}(A_{m-1}B_mO_{3m+1})^{2-}$ (where A is a mono-, di-, or trivalent cation with 12-fold-coordination, B is a transition-metal cation with 6-fold-coordination and m is the number of octahedral layers in the perovskite slab)[6] with the addition of magnetic species on the B-site have been investigated and showed to be one of the new leading alternatives to $BiFeO_3$ as a single-phase multiferroic material.[7-11] The Aurivillius phase materials $Bi_5FeTi_3O_{15}$ (m=4) and $Bi_6Fe_2Ti_3O_{18}$ (m=5) which can be considered as a solid solution of $Bi_4Ti_3O_{12}$ and $BiFeO_3$, showed coexistence of FE and FM order.[12,13] By substituting Co and Fe in the B-site, the ferromagnetism of $Bi_5Fe_{0.5}Co_{0.5}Ti_3O_{15}$ (m=4) and $Bi_6FeCoTi_3O_{18}$ (m=5) was remarkably improved with both compositions showing multiferroic properties at room temperature.[14,15] However, the origin of their magnetic property is still under debate because the net magnetization may originate from the presence of secondary phases which are not easily detected by laboratory XRD diffractometers.[16] Therefore, more research was needed to further understand the origins of the multiferroic properties of Aurivillius phase materials.

In this study, textured Aurivillius phase $Bi_{5.25}La_{0.75}FeCoTi_3O_{18}$ (BLFCT) ceramics were prepared by Spark Plasma Sintering (SPS). Partial La substitution was chosen to reduce the electrical conductivity and to decrease the ferroelectric coercive field, while

Co, in co-substitution with Fe was selected to increase magnetization.[14,17] Due to the structural anisotropy of Aurivillius phase materials and the presence of ferroelectric polarization in the *a-b* plane, the SPS textured BLFCT ceramic, would allow easier switching of the polarization in the *a-b* plane if compared to that prepared by conventional sintering. The ferroelectric behaviour of the Aurivillius phase is confirmed by the presence of current peaks in the current vs. electric-field loops which indicate the occurrence of ferroelectric domain switching.[18] The ferromagnetic character was indicated by the magnetic field-induced reversible intensity changes of a certain set of crystalline planes belonging to the Aurivillius phase, measured by in-situ neutron diffraction under applied magnetic field on the WISH instrument at the ISIS Neutron Spallation facility, RAL, Didcot, Oxfordshire, UK. Ferroelectric domain switching induced by an applied external magnetic field suggests the existence of spontaneous toroidal moments and a linear magnetoelectric coupling in the material at room temperature (RT). Therefore, we can conclude that the textured Aurivillius BLFCT phase material prepared by SPS is an intrinsic multiferroic with the coupled ferroelectric and magnetic order parameters at room temperature.

The BLFCT samples were obtained using the conventional solid state reaction route and SPS. The details of sample preparation, characterization and first principles modeling of magnetization are showed in supplementary information.

Figure 1 shows the X-ray diffraction (XRD) pattern of the BLFCT powder. All the peaks were indexed to an orthorhombic structure with a *B2cb* space group characteristic of five-layered Aurivillius phase compounds, and simultaneously to another orthorhombic phase with a crystallographic symmetry of a $A2_1am$ space group, typical of four-layered Aurivillius oxides.[19,20] The lattice parameters of the orthorhombic *B2cb* structure $a = 5.4530(2)$ Å, $b = 5.4427(1)$ Å and $c = 50.670(2)$ Å, and the four-layered $A2_1am$ phase $a = 5.4651(6)$ Å, $b = 5.3943(6)$ Å and $c = 41.487(2)$

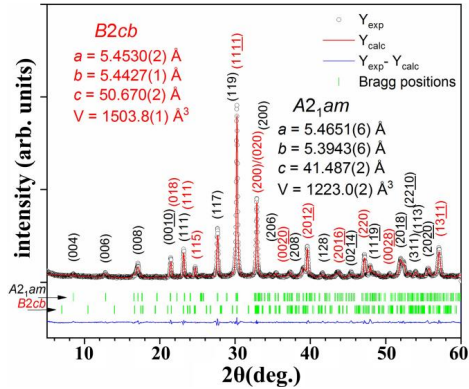Å were obtained by the full pattern matching method using the FullProf software package (http://www.ill.eu).

FIG. 1. X-ray diffraction pattern of the BLFCT powder.

These results indicate that BLFCT is a mixed-layer Aurivillius material consisting of the m=4 and m=5 perovskite slabs. No extra peaks of possible impurity phases were identified within the XRD's instrumental resolution. Fig. S1 shows the SEM images of the textured BLFCT ceramics. The plate-like grains (*a-b* axes) are aligned in the direction perpendicular to the SPS pressure direction. Moreover, a small amount of a secondary phase also marked by the red circle in Fig. S1 can be observed in the microstructure. The amount of this secondary phase is estimated to be approximately 1.4 vol. %, which is below the trace level of XRD. EDX line scanning analysis (Fig. S2 and Table S1) revealed that, local chemical inhomogeneities are composed of Fe, Co, O and Ti. Thus, the secondary phase is assumed to be of the Co_2FeO_4 spinel ferrite type with cobalt ions partially replaced by titanium ions. Such a chemical impurity has also been reported to exist in four-layer Aurivillius compounds $Bi_5Fe_{0.5}Co_{0.5}Ti_3O_{15}$.[16]

Figure 2a shows the temperature dependence of dielectric permittivity and loss of BLFCT ceramics measured perpendicular to the SPS pressure direction at five different frequencies (50, 70 100, 300 and 500 kHz). The peak of the dielectric permittivity at

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about 1060 K corresponds to the FE T_c of BLFCT. Hence, the ferroelectric transition temperature of BLFCT is higher than that of $Bi_6Ti_3Fe_2O_{18}$ phase (at 973 K).[13] Fig. 2b shows the P-E and I-E hysteresis loops of the textured BLFCT ceramics measured in a perpendicular direction to the SPS pressure at RT. The observed current peaks in the I-E loops confirm the ferroelectric domain switching, [21,22] which occurs in the sample upon field reversal. The maximum value of the polarization of BLFCT is 10 μC/cm^2.

Figure 2c shows the temperature dependence of the zero-field cooled (ZFC) and then field-cooled (FC) magnetization under an applied magnetic field of 200 Oe for BLFCT. We propose that BLFCT undergoes a paramagnetic to ferromagnetic transition at T_{c1} ~ 494 K (defined as the temperature corresponding to the peak of dM/dT) which is close to that of $Bi_6FeCoTi_3O_{18}$ (526 K).[23] The ferromagnetism of BLFCT is induced by the double exchange interactions between neighbouring Fe^{3+}-O-Fe^{3+}, Co^{3+}-O-Co^{3+} and Fe^{3+}-O-Co^{3+} (cf. the following density functional theory calculations).[24] This is consistent with the EDX result for the elemental make-up of the main phase. Another peak clearly observable on the ZFC curve at T_{c2} ~ 353 K can be associated with the magnetic transition of the secondary phase Ti doped Co_2FeO_4. The fact that T_{c2} is lower than the transition temperature of Co_2FeO_4 (460K) can be explained by Ti substitution.[16, 25] The spontaneous magnetization (M_s) of Co_2FeO_4 is about 16 - 23.5 emu/g.[25] Therefore, spontaneous magnetization of 1.4 vol.% $Co_{2-x}FeTi_xO_4$ should be about 0.22 - 0.32 emu/g which is much smaller than the spontaneous magnetization of BLFCT (M_s = 1.85 emu/g, Fig.2d). In order to distinguish the contribution of secondary phase and main phase, the M-H loop was measured at a temperature between the two phase transition T_{c1} and T_{c2} (**Figure S3**). The spontaneous magnetization of the ceramic measured at 425 K is about 1.58 emu/g. This means that the magnetic contribution of the secondary phase is about 0.27 emu/g which agrees with the secondary phase amount obtained from the EDX results.

Therefore, the main ferromagnetism of BLFCT is believed to originate from the Aurivillius phase.
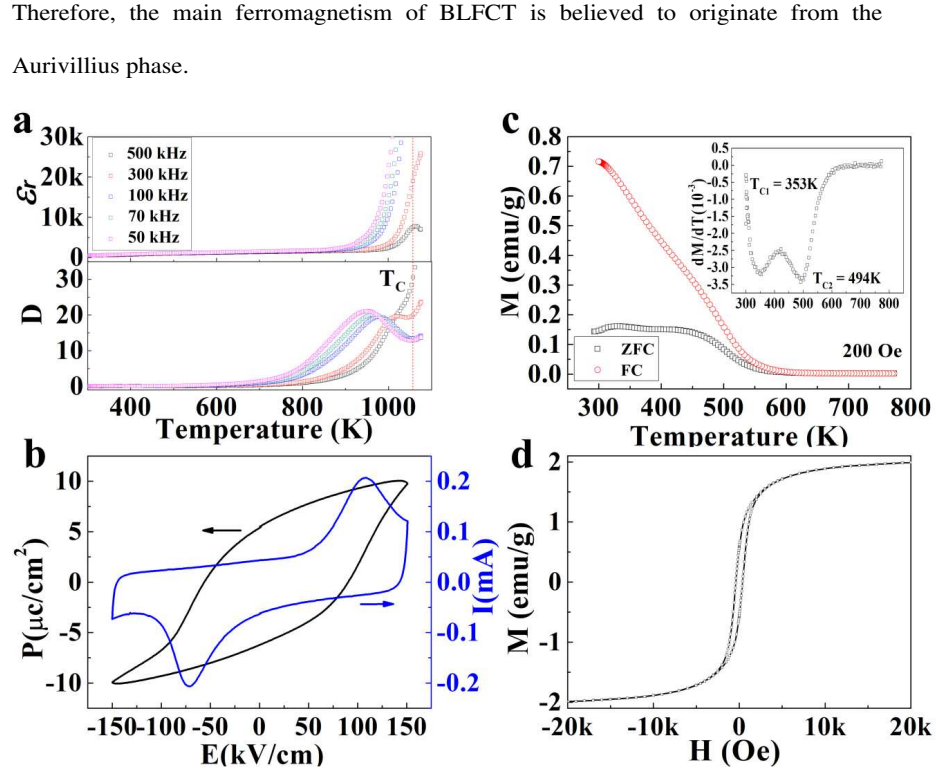

FIG. 2. (a) The temperature dependence of the relative dielectric permittivity (above) and loss tangent (down) of BLFCT ceramics measured perpendicular [⊥] to the SPS pressure direction; (b) The P-E and I-E hysteresis loops of the textured BLFCT ceramics measured perpendicular to the SPS pressure direction. The curves demonstrate the domain switching characteristic of a ferroelectric. (c) The temperature dependence of the magnetization for BLFCT ceramics [⊥] in the ZFC/FC modes measured in H = 200 Oe. The inset shows the first derivative of magnetization dM/dT as a function of temperature; (d) The field dependence of magnetization at 300K for the BLFCT sample.

Figure 3 demonstrates an example of the ferromagnetic exchange mechanism between neighbouring Fe^{3+}-O-Co^{3+} by the density functional theory (DFT) calculations with the Vienna ab initio simulation package (VASP) code. Effective Hubbard interaction of U_{Fe} = 2 eV and U_{Co} = 3 eV for Fe and Co atoms, respectively, were

introduced to take strong correlation into account in the generalized gradient approximation (GGA)+U method. It shows that the magnetism of the BLFCT is mainly originated from the transition metal ions (Fig.4a), Fe^{3+} and Co^{3+} (3.1 and 2.1 $\overline{\mu_B}$/atom, respectively), while their surrounding O and Ti atoms have a small contribution as well (less than $0.1\overline{\mu_B}$/atom). The exchange processes between the adjacent corner-shared FeO_6 and CoO_6 tilted octahedra are dominated by the (virtual) electron transfer between the Fe/Co spin-up d_{xy} orbitals at the Fermi level through the O spin-up p_x/p_y orbitals (cf. Fig.3b). For characterizing the type of magnetic coupling, two different relative alignments of the magnetic moments of Fe^{3+} and Co^{3+} cations, parallel (i.e., ferromagnetic) alignment and antiparallel (i.e., antiferromagnetic) alignment, are both considered. The ferromagnetic orientation is found to have much lower energy than the antiferromagnetic orientation and a large energy difference reads E_{FM}-E_{AFM} = -144.1 meV. The double exchange ferromagnetism occurs in this doped system of localized spins and itinerant electrons interacting via the Hund's rule coupling. The strength of the (FM) exchange interaction is around of 43.5 meV (i.e. 504.6 K), which is consistent with the FM transition temperature T_{c1} in the ZFC/FC measurements (cf. Fig. 2c). The parallel alignment of magnetic moments is identified within the *a-b* plane. Together with the following neutron diffraction measurements, the in-plane orientation of magnetic moments is found to be along the [010] direction.
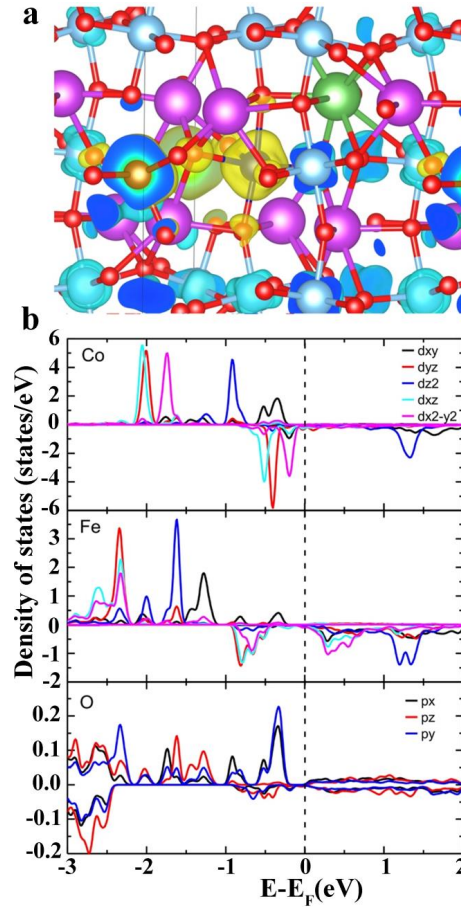
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FIG. 3. (a) Spin density of the system (Isosurface level = 0.005): O (red), Bi (pink), La (green), Ti (light blue), Co (dark blue) and Fe (yellow) atoms. (b) A detailed inspection of density of state in Fe^{3+}-O-Co^{3+}.

Figure 4 shows the neutron diffraction data collected for the BLFCT sample subjected to an external magnetic field. Intensity changes in the diffraction pattern can be observed in several peaks at large d-spacing, while at low d-spacing no observable intensity change is detectable. This observation is in agreement with the expected dependence from the magnetic form factor. No additional peaks could be detected, indicating that the magnetic structure is commensurate with the crystallographic structure, with possible ferromagnetic or weak

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ferromagnetic character. In the inset of Fig. 4, it is shown that the intensity of the significant Aurivillius peak (0 1 20) reversibly changes under the full excursion of the applied magnetic field. The peaks' intensity increases with increasing field up to 2T. For higher field ($2T < H <$ 5T), the peak intensity remains unchanged, which indicates that the magnetic material becomes saturated at about 2T, in agreement with the M-H loops in Figs. 2d and S3. Finally, a remanent magnetic contribution is observable when compared to the zero field measurements.
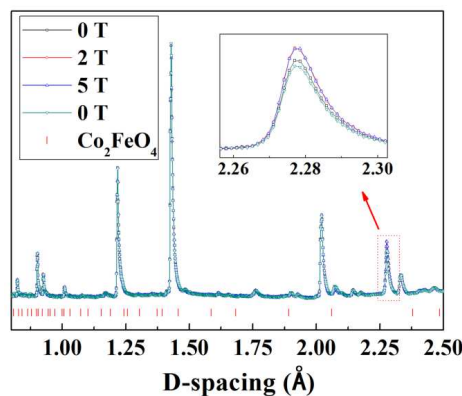
FIG. 4. The room-temperature neutron diffraction pattern obtained for BLFCT ceramic at different magnetic fields (from 0, 2, and 5 T then back to 0T).

Figure 5 demonstrates the existence of magnetoelectric coupling in the BLFCT sample by the PFM images of changes in ferroelectric domain structure induced by an applied magnetic field greater than the magnetic coercive field 399 Oe. Vertical PFM was performed on the BLFCT sample as shown in Fig. 5d. A partial ferroelectric domain switching observed in the local Fe and Co- rich nanostoichiometric regions (marked by the red circles) was obtained after applying an in-plane magnetic field of +2000 Oe (Fig. 5e). When a magnetic field of - 2000 Oe is applied in the opposite direction, part of the ferroelectric domains switch back (Fig. 5f). The low content of the cubic (non-polar) secondary phase Ti doped Co_2FeO_4 and its non-ferroelectric character exclude the possibility that the described switching phenomena might be due to the secondary phase. Therefore, the magnetic-field driven changes in the

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ferroelectric domain structure are intrinsic, resulting from the magnetoelectric response coupling in the SPS prepared Aurivillus BLFCT ceramic. Given the fact that the ferroelectric polarization ($\overline{P}$) lies along the [100] direction and the ferromagnetic moments ($\overline{M}$) are mostly around the [010] direction, the presence of a spontaneous (ferro-)toroidal moment $\boxed{T = P \times M}$ and a linear magnetoelectric effect is suggested to occur in BLFCT ceramics.
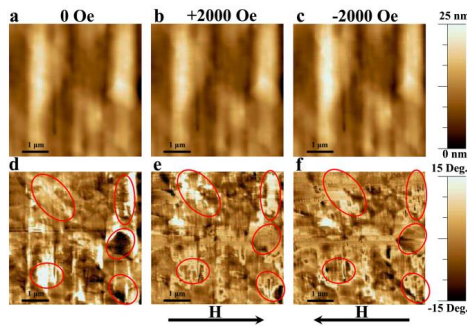

FIG. 5. Ferroelectric domain switching in BLFCT under applied magnetic field (H is parallel to the surface of sample), (a) Topography and (d) vertical PFM phase at 0 Oe; (b) topography and (e) vertical PFM phase at +2000 Oe; (c) topography and (f) vertical PFM phase at -2000 Oe.

In conclusion, the Aurivillius phase in BLFCT ceramic is found to be ferroelectrically and ferromagnetically active at room temperature. Ferromagnetism is suggested to originate from double exchange interactions between neighbouring Fe^{3+}-O-Fe^{3+}, Co^{3+}-O-Co^{3+} and Fe^{3+}-O-Co^{3+} in the main phase by the first principle calculations. Magnetic controlled ferroelectric domain switching was observed by piezoelectric force microscopy at room temperature. The Aurivillius phase ceramics, with Co/Fe contributing to magnetization and polarization at the same time, can be considered an intrinsic room-temperature multiferroic with ferro-toroidal phase.

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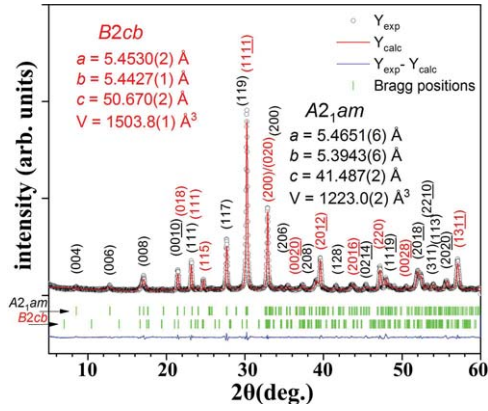
DATA AVAILABILITY STATEMENT

The data that supports the findings of this study are available within the article [and its supplementary material].
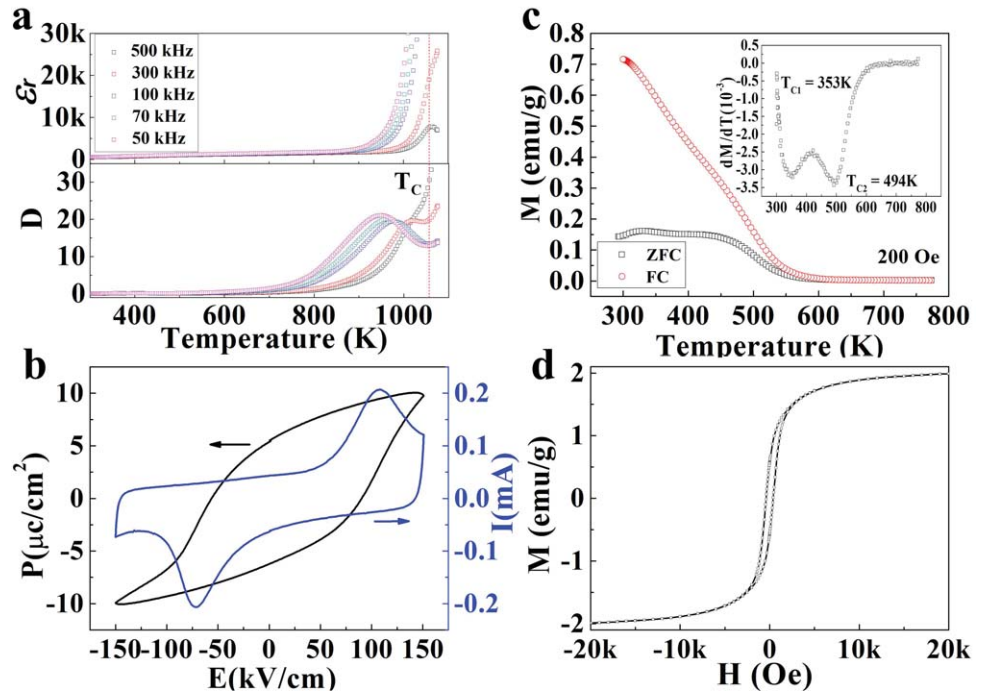
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